

1.3 µm-Wavelength Laser with Type-II Active Layer on GaAs Substrate

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density integration of optical devices. The challenge lies in the potential degradation of device performance due to heat To meet the rapidly increasing demand for optical communication traffic, a high-capacity solution is required through the highgeneration. Consequently, there is a need for a semiconductor laser exhibiting superior temperature attributes. Addressing this, we have developed a 1.3 µm-wavelength laser, employing a GalnAs/GaAsSb/GalnAs type-II active layer on a GaAs substrate, suitable for optical communication. This has been achieved utilizing our proprietary low-temperature growth technology from highly productive Organic-Metal Vapor Phase Epitaxy. Our solution exhibits a characteristic temperature of the threshold current density of 152 K between 25°C and 100°C, a value significantly larger than the traditional 60 K of InGaAsP lasers on InP substrates. This result demonstrates the feasibility of a 1.3 um semiconductor laser with a type-II active layer. Jess susceptible to performance degradation in high-temperature environments.

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Keywords: GalnAs, GaAsSb, GaAs, type-II, OMVPE

Introduction 1.

While various applications provided by advanced information and communication technologies, such as generative AI, cloud services, and automated driving, are improving convenience in our daily lives, communication traffic is rapidly increasing. Therefore, network switches in data centers are required to simultaneously increase tion architectures such as Co-Packaged Optics (CPO) and capacity and reduce power consumption, and next-generaoptical devices combining III-V compound semiconductors and silicon (Si) photonics have been proposed as solutions.

grated to achieve large capacitance, but the degradation of In these proposals, optical devices are densely intecharacteristics due to heat generation from the optical anism using a Peltier element is employed to address the devices is a problem. Usually, a temperature control mechheat generation problem, but this has disadvantages in terms of integration, power consumption, and cost. ductor lasers with excellent temperature characteristics that Therefore, there is a need for Peltier element-free semicondo not easily degrade even in high-temperature environ-
ments.

Semiconductor lasers, which are primarily used in optical information communications, operate in the $1.3 \mu m$ band, where fiber dispersion is minimal, and in the 1.55 µm band, where losses are minimal. Semiconductor lasers in these wavelength bands generally use InGaAsP-based materials on InP substrates, but they mainly suffer from teristics such as increased threshold current and decreased flow, $^{(2)}$ and there is significant degradation of laser characissues such as Auger recombination⁽¹⁾ and carrier overoptical output power with temperature rise.

nation process in which an electron and a hole recombine Auger recombination is a non-luminescent recombito excite other carriers without emitting light, which is converted into thermal energy. The probability of Auger recombination increases as the band gap energy of the semiconductor where the optical transition occurs

decreases. In addition, the higher the carrier concentration, the more pronounced the Auger recombination is, which further increases the rise in the threshold current due to temperature rise and leads to degradation of the tempera-
ture characteristics.

Career overflow is a phenomenon in which electrons leak out of the active layer and no longer contribute to light emission recombination due to thermal energy. In materials suitable for the growth of InGaAsP-based materials on InP substrates with lattice-matched semiconductor materials. there is no material with a large band gap, and electron leakage becomes significant due to the small band offset between the active layer and the clad layer. On the other hand, on GaAs substrates, lattice-matched semiconductor materials with a large band gap can be used, enabling the ture characteristics. However, in GaInAs-based materials realization of semiconductor lasers with excellent temperasion wavelength by increasing the In composition. used for the active layer, it is possible to increase the emis-Nevertheless, as the In composition increases, the crystal quality deteriorates significantly due to lattice strain with the GaAs substrate, making it extremely difficult to achieve long wavelength extension up to the wavelength band for optical communication.

Therefore, we focused on the GaInAs/GaAsSb system, which can be extended to longer wavelengths with low In composition and suppressed Auger recombination by using type-II transitions on GaAs substrates. We have succeeded in fabricating GaInAs/GaAsSb superlattice on InP substrates for near-infrared sensors by the Organic-Metal ized a high-performance near-infrared sensor. (3) , (4) lent mass-producibility, and have realized and commercial-Vapor Phase Epitaxy (OMVPE) method, which has excel-In this study, we report on the development of the growth technology of GaInAs/GaAsSb type-II active layers on GaAs substrates by applying our original low-temperature OMVPE growth technology cultivated in the development of these lasers, and on the fabrication of type-II active-layer

ridge-waveguide lasers on a 1.3 μm-band GaAs substrate suitable for optical communications, that is Fabry-Perot (FP) lasers.

2. Type-II Quantum Well Structure and Growth

2-1 Type-II quantum well structure

conductors with different band gaps stacked on the order of First, a quantum well (QW) structure consists of semiseveral nm. Unlike the bulk structure, the optical gain is greatly improved, making it an indispensable structure for conductor lasers used in the field of optical information the active layer of semiconductor lasers. Almost all semiand telecommunications employ the type-I structure, in which electrons and holes are confined in the same well layer and optical transitions occur in the active layer. On the other hand, there are type-II quantum wells, such as the GaInAs/GaAsSb system shown in Fig. 1, in which the rated and optical transitions occur due to the overlap of the wavefunctions of electrons and holes are spatially sepawavefunctions near the hetero-interface. Advantages include the following:

- tive transition energy while using semiconductors (1) Enables longer wavelengths by reducing the effecwith a large band gap
- (2) Suppression of Auger recombination⁽⁵⁾

When utilizing this type-II transition in the active layer of semiconductor lasers, a GaInAs/GaAsSb/GaInAs structure (so-called "W" type), in which a hole-trapping GaAsSb layer is sandwiched by an electron-trapping GaInAs layer, is used. The overlap of the wavefunctions istics by increasing the recombination probability. $(6)(7)$ increases, which is expected to improve the laser character-

Fig. 1. Band structure of GaInAs/GaAsSb type-II quantum wells

2-2 Growth of type-II quantum well structure

The OMVPE method, known for its high productivity, was employed as the growth technique. Firstly, to verify the effect of long-wavelength extension through type-II transitions, $GaAs/GaIn_{0.34}As-QW$ with a thickness of 4 nm, $GaAs/GaAsSb_{0.23}-QW$, were combined to form $GaIn_{0.34}As/$ $GaAsSb_{0.23}/GaIn_{0.34}As "W" - QW. Additionally, for further$ $long-wavelength$ extension, $Galn_{0.34}As/GaAsSb_{0.30}/$ $Galn_{0.34}As''W''-QW$ was grown by increasing the Sb.

composition. The band structure and photoluminescence (PL) measurements at room temperature are shown in Fig. 2. 1030 nm wavelength for $GaAs/GaIn_{0.34}As-QW$, 1130 nm for $GaAs/GaAsSb_{0.23}-QW$ and $GaIn_{0.34}$ As/ $GaAsSb_{0.23}/GaIn_{0.34}As "W"-QW emits at a wavelength of$ 1230 nm, confirming the effect of longer wavelengths due to type-II transitions, as expected, and $Galn_{0.34}As/$ $GaAsSb_{0.30}/GaIn_{0.34}As "W"-QW emission at 1360 nm was$.achieved

Fig. 2. Band structure and PL spectra of GaAs/GaInAs-QW, GaAs/GaAsSb-QW, and GaInAs/GaAsSb/GaInAs "W"-QW at room temperature

ro-interfaces, so the quality of the interface is very As mentioned above, type-II transitions occur at hetero-interfaces is difficult due to the effects of elemental important. Generally, the formation of element-rich heteinterdiffusion and segregation, which make the formation of steep interfaces difficult. In this study, we attempted to perature OMVPE growth technique. Figure 3 shows PL improve the interface quality using our original low-temspectra before and after the improvement of interface quality. Specifically, by lowering the growth temperature (T_{sub}) from 550°C to 475°C, the PL intensity was improved by a factor of three, and the Full Width at Half Maximum $(FWHM)$ was improved from 88 nm to 75 nm. Figure 4 shows a cross-sectional transmission electron microscope (TEM) image of the GaInAs/GaAsSb hetero-interface at ro-interface. Since the contrast of this intermediate layer is ence of an intermediate layer at the GaInAs/GaAsSb hete- 550° C. The black contrast in the figure indicates the presclose to that of the GaAs layer, it is assumed to be a component with atomic weight similar to that of GaAs, and it is inferred that In atoms segregate in the initial growth stage of the GaInAs layer and Sb atoms segregate in the initial growth stage of the GaAsSb layer. On the other hand, the formation of the intermediate layer was suppressed in the 475° C growth, suggesting that the low-temperature growth suppressed the segregation of In and Sb and improved the steepness of the hetero-interface, which contributed to the improvement of luminescence properties.

Fig. 3. PL spectra at growth temperatures of 550° C and 475° C

Fig. 4. Cross-sectional TEM images at growth temperatures of 550° C and 475° C

3. Fabrication of Ridge-waveguide FP Lasers

Figure 5 shows the cross-sectional structure of the fabricated ridge-waveguide FP laser. The structure was grown on an n-type GaAs substrate using OMVPE method, which consists of an $n - Al_{0.45}GaAs$ cladding layer (1500) nm), GaAs SCH layer (150 nm), GaInAs (4 nm)/GaAs Sb (4 nm) GaInAs (4 nm) "W"-QW active layer with 2 pairs, a GaAs SCH layer (150 nm), p-Al0.45GaAs cladding layer (1500 nm) , and p-GaAs contact layer (200 nm) . Subsequently, a multi-level ridge structure was fabricated with a mesa width of $2-30 \mu m$ using common photolithog-

Fig. 5. Structure of the fabricated ridge-waveguide FP laser

wafer was passivated with an insulating film, and an ohmic electrode was deposited by vacuum evaporation with the insulating film opened only directly above the p-GaAs contact layer. The backside of the wafer was polished to reduce the thickness to approximately $100 \mu m$, and then the backside electrode was formed. The end face was guide FP lasers with resonator lengths of 1-3 mm were formed by cleavage, and multiple levels of ridge-wavefabricated. Finally, the element was epi-up mounted in a

Characteristics Laser 4.

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raphy and we technique centire surface of the missional wether be control and the insulating film, and an ohmic electrode was deposited by vacuum evaporation with the insulating film opened only directly above the p-GaAs Figure 6 (a) shows the current density-optical output characteristics of the device with a mesa width of 5 µm and a resonator length of 1 to 3 mm at room temperature and pulse driving (pulse width: 10 usec, duty ratio: 1%), and (b) shows the oscillation spectrum. For the 1 mm resonator length device, a threshold current density of about 1.8 $kA/cm²$ and laser oscillation at a wavelength of 1286 nm dence of the threshold current in a device with a resonator were achieved. The evaluation of the mesa width depenlength of 1 mm confirms the absence of reactive current flowing beyond the active layer. From the resonator length ciency, the transparency current density per "W"-QW is dependence of the threshold current density and slope effiestimated to be 155 A/cm² and the mode gain is 12.7 cm⁻¹, which are comparable to those of InP-based materials. On the other hand, the internal quantum efficiency^{*1} is esti mated to be 33% , which is lower than that of a typical ciency is lower for type-II quantum wells because the type-I laser. It is assumed that the internal quantum effioverlap of wavefunctions is reduced compared to type-I. In order to isolate whether the obtained values are due to intrinsic or process-induced factors of the type-II active layer, we fabricated a gain-guided FP laser device without a ridge structure with an electrode width of 13 um.

Fig. 6. (a) Current density-optical output characteristics of an element with a mesa width of 5 μ m and a resonator length of 1 to 3 mm at room temperature and pulse driving (pulse width: 10 µsec, duty ratio: 1%) (b) Oscillation spectrum

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Consequently, an internal quantum efficiency of 58% was nation efficiency prompted by the ridge waveguide sides, attained, attributed to a decline in luminescence recombinecessitating enhancement.

characteristics Temperature 4-2

Figure 7 (a) shows the current density-optical output of a 5 μ m mesa width, 1 mm resonator length device at $25-100^{\circ}$ C in pulsed operation (pulse width 10 µsec, duty ratio 1%). Clear laser oscillation was achieved at each temperature. The characteristic temperature^{*2} T_0 for the threshold current density was 333 K at 25-50°C, 164 K at 50-75°C, 92 K at 75-100°C, and 152 K at 25-100°C (Fig. 7 (b)), while the existing InGaAsP-based laser on an InP substrate has a characteristic temperature of about 60 $K^{(8)}$ It has also been reported that the characteristic temperature increases with a smaller threshold current

Fig. 7. (a) Current density-optical output at $25{\text -}100^{\circ}\text{C}$ for the pulsed drive (pulse width 10 µsec, Duty ratio 1%) of a 5 μ m mesa width and 1 mm resonator length device

(b) Temperature dependence of threshold current density

Fig. 8. (a) Oscillation spectra of an element with a mesa width of 5 μ m and a resonator length of 1 mm at 25-100°C

(b) Temperature dependence of the oscillation wavelength

density per unit quantum well layer due to the Auger effect. (9) Considering that the threshold current density per based lasers due to the lower internal quantum efficiency "W"-QW at 25° C is slightly higher than that of InGaAsPcompared to type- $I₁⁽⁹⁾$ the temperature characteristics of type-II active layer lasers on GaAs substrates are extremely .good

The oscillation spectrum at $25{\text -}100^{\circ}\text{C}$ is shown in Fig. 8 (a). For the temperature dependence coefficient of the oscillation wavelength d λ /dT, we obtained 0.43 nm/K for 25-50°C, 0.29 nm/K for 50-75°C, and 0.25 nm/K for cient of the oscillation wavelength became lower than that 75-100 $^{\circ}$ C (Fig. 8 (b)). The temperature dependence coeffiof the existing InGaAsP-based lasers $(0.4$ to 0.5 nm/K) in the temperature range above 50° C.⁽⁸⁾ The temperature dependence of the oscillation wavelength is the subtraction of the gain peak lengthening effect with increasing temperature and the shortening effect due to the band bending effect caused by the injected carriers. The higher the temperature dependence of the injection carrier amount required for laser oscillation, i.e., as the characteristic temperature decreases, the shortening effect becomes stronger and the temperature dependence coefficient of the oscillation wavelength decreases as well.⁽⁸⁾ Compared to icant band bending effects in relation to the injection type-I quantum wells, type-II quantum wells exhibit signifcarrier amount.^{(10)} It is assumed that type-II active layer lasers on GaAs substrates achieved a low temperature dependence coefficient of oscillation wavelength while ately designing this band bending effect, it is expected that maintaining a high characteristic temperature. By approprisemiconductor lasers with high characteristic temperature and temperature-independent oscillation wavelength can be .realized

Conclusion 5.

In this study, we fabricated ridge-waveguide FP lasers by growing a GaInAs/GaAsSb/GaInAs type-II active layer etary low-temperature OMVPE growth technique. Laser on a GaAs substrate through the utilization of our propritures ranging from 25° C to 100° C, the characteristic oscillation at 1.3 μm band was achieved, and at temperatemperature of the threshold current density significantly exceeded 152 K, surpassing the value of 60 K obtained with conventional InGaAsP lasers on InP substrates. By using GaInAs/GaAsSb/GaInAs type-II active layer on ductor lasers that are less prone to degradation of laser GaAs substrates, it is possible to achieve $1.3 \mu m$ semiconcharacteristics even in high-temperature environments.

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Technical Terms

- *1 Internal quantum efficiency: Percentage of carriers injected from electrodes that are converted to light.
- *2 Characteristic temperature: The threshold current density of the laser, J_{th} , is given for the active layer *T* temperature

 $J_{\text{th}} \propto \exp(T/T_0)$

 T_0 at this time is called the characteristic temperature. The higher the characteristic temperature, the more difficult it is for the laser characteristics to deteriorate even at higher.

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